

# **Study of Top Quark Rare Decays Via FCNC at LHC**

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# Introduction

- LHC can be considered as a “**top factory**”, producing about  $10^8$   $tt^-$  pair per year at luminosity  $L=100 \text{ fb}^{-1}$ .
- The decays (**FCNC**)  $t \rightarrow Zq$  ( $q=u,c$ ),  $t \rightarrow Hq$  ( $q=u,c$ ) is strongly suppressed (**loop suppression, heaviness of gauge bosons**) in **SM**. The **SM** predicts :  
 $Br(t \rightarrow Zq) \sim 1.3 \times 10^{-13}$   
 $Br(t \rightarrow Hq) \sim 0.9 \times 10^{-13} (4 \times 10^{-15})$  for  $m_H = 100(160) \text{ GeV}$
- Observation of significantly larger Br's will be a **clear signal of new physics** (**new dynamical interactions of top quark, multi-Higgs doublets, exotic fermions or other possibilities**).
- The recent theoretical (MSSM, 2HDM II, Technicolor TC2) estimations:  
 $Br(t \rightarrow Zq) \sim 10^{-5} \div 10^{-4}$   
(J.Yang, B..L.Young and X.Zhang, Phys Rev., D58 (1998) 055001)  
 $Br(t \rightarrow Hq) \leq 5 \times 10^{-4}$   
(J.Guasch and J.Sola, Nucl. Phys., B562 (1999) 3: hep/ph/9906268, 1999)
- The analysis presented here focus on the following final state topology of  
 $tt^- \rightarrow HqWb \rightarrow W W^* q W b \rightarrow l\nu l\nu j, l^\pm \nu b$  ( $l=e,\mu$ ).

# FCNC Decay $t \rightarrow Zq$ ( $q=u,c$ )

- The sensitivity of the **ATLAS experiment** to the  $t \rightarrow Zq$  decay has been studied in two decay modes:

$tt^- \rightarrow ZqWb \rightarrow l^+ l^- j, l^\pm \nu b$  ( $l=e, \mu$ ) – leptonic final state topology

$tt^- \rightarrow ZqWb \rightarrow l^+ l^- j, jjb$  ( $l=e, \mu$ ) --- hadronic final state topology

For  $t \rightarrow Zq$  decay the following SM backgrounds has been considered:

- $Z+jets \rightarrow l^+ l^- + jets$
- $tt^- \rightarrow WbWb \rightarrow l^+ \nu b l^- \nu b$
- $pp \rightarrow WZ + X \rightarrow l^\pm \nu l^+ l^- + X$

It has been obtained that, the  $Br(t \rightarrow Zq \rightarrow l^+ l^- q)$  as low as  $2.0 \times 10^{-4}$  for the **leptonic mode** and as low as  $5.9 \times 10^{-4}$  for the **hadronic mode** could be discovered at the  $5\sigma$  level with an integrated luminosity of  $100 \text{ fb}^{-1}$ .

# Monte Carlo Events Generation

The signal  $t \rightarrow Hq$  has been implemented in the **PYTHIA 5.7**. The following SM backgrounds has been considered:

- $tt^- \rightarrow W^+ b W^- b \rightarrow l^+ \nu b l^- \nu b$
- $tt^- H \rightarrow W^+ b W^- b, WW^* \rightarrow l^+ \nu b, l^- \nu^- b^-, l^+ \nu l^- \nu^-; l^\pm \nu b, jjb, l^+ \nu l^- \nu^-; l^+ \nu b, l^- \nu^- b^-, jj l \nu$
- $WZ \rightarrow l^\pm \nu l^+ l^- + X$
- $WH \rightarrow l^\pm \nu WW^* \rightarrow l^\pm \nu, l^+ \nu l^- \nu^- + X$

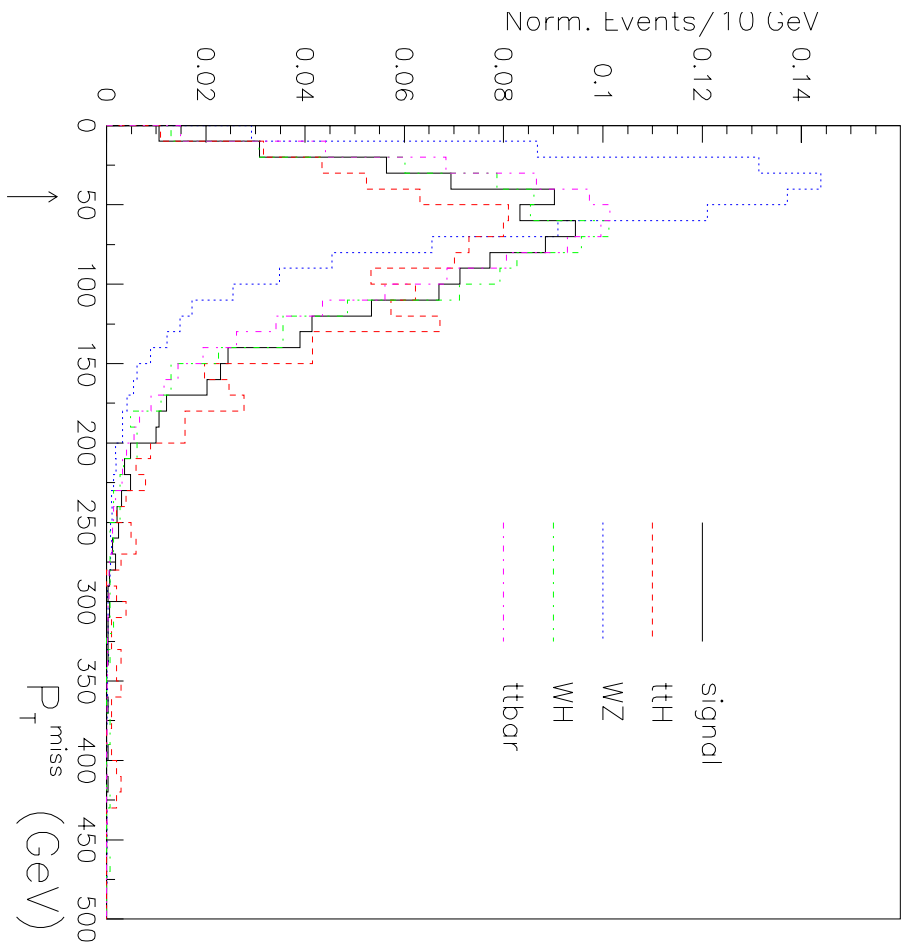
- All backgrounds were generated by **PYTHIA 5.7** at  $\sqrt{s} = 14 \text{ TeV}$ ,  $m_{\text{top}} = 175 \text{ GeV}$  for two masses of Higgs  $m_H = 150 \text{ GeV}$  and  $m_H = 160 \text{ GeV}$ , with proton structure function **CTEQ2L**. Initial and final state QED and QCD (ISR, FSR) radiation, multiple interactions, fragmentations and decays of unstabled particles were enabled. The cross-section for  $tt^-$  production was assumed to be  $\sigma_{tt} = 833 \text{ pb}$ .
- The events were simulated using the detector fast simulation package **ATLFAST 2.51**. The b-tagging performance was simulated assuming the nominal efficiencies of  $\epsilon_b = 50 \%$ ,  $\epsilon_c = 10 \%$ ,  $\epsilon_j = 1 \%$ .
- The branching ratios of the  $H \rightarrow WW^*$  decays have been estimated by the Fortran code **HDECAY** (a. Djouadi, J.Kalinovski and M.Spira, hep-ph/9704448,1997)  $\text{Br}(H \rightarrow WW^*) = 0.68$  for  $m_H = 150 \text{ GeV}$  and  $\text{Br}(H \rightarrow WW^*) = 0.92$  for  $m_H = 160 \text{ GeV}$ .

# Event Analysis

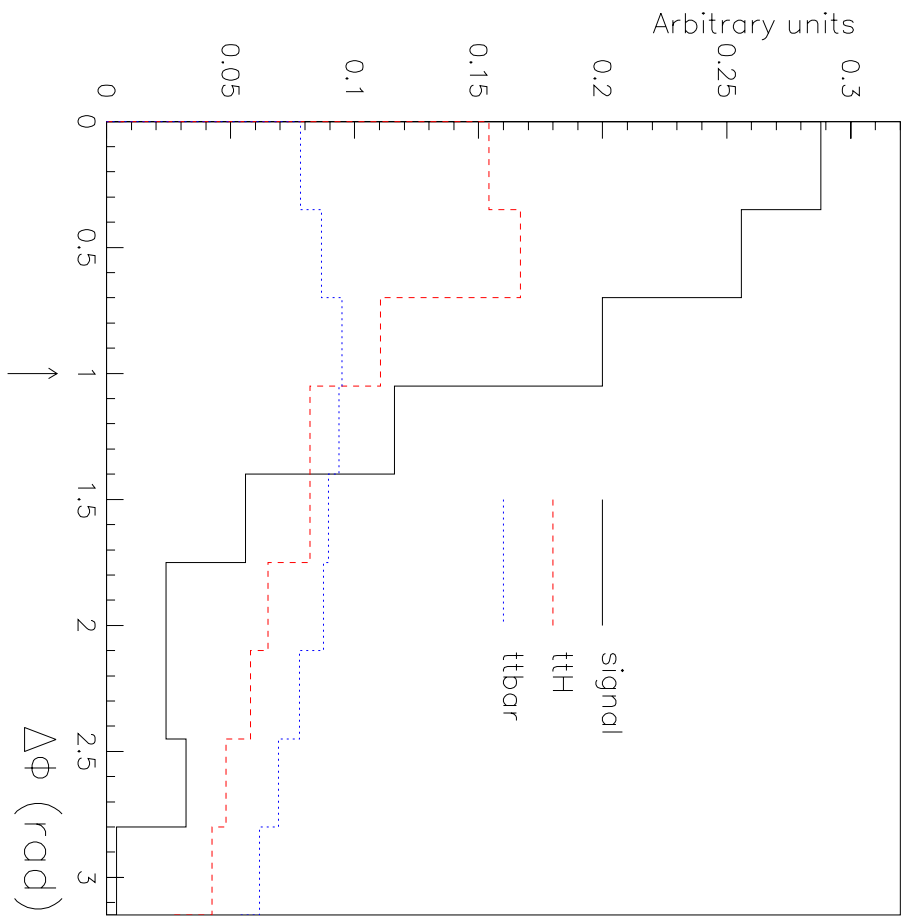
- **Preselection cuts: the presence of at least 3 charged leptons (electron with  $P_T > 5 \text{ GeV}$ ,  $|\eta| < 2.5$  and muon with  $P_T > 6 \text{ GeV}$ ,  $|\eta| < 2.4$ ); the number of jets with  $P_T^{\text{jet}} > 15 \text{ GeV}$ ,  $|\eta^{\text{jet}}| < 5$  at least 2. The preselection cuts reduce backgrounds up to  $2 \div 14\%$ , meanwhile  $\sim 50\%$  of signal events are survived.**
- **The requirement of 3 isolated leptons with high  $P_T^l > 30 \text{ GeV}$  certainly includes the presence of one opposite charged lepton. This cut affects significantly  $t\bar{t}$  dangerous background, while keeping still a significant fraction of the signal.**
- **The next requirement of the missing transverse momentum in event with  $P_T^{\text{miss}} > 45 \text{ GeV}$  is a powerful cut for reduction of  $WZ$  background, though other backgrounds are less sensitive to this cut.**
- **We demand the presence of at least 2 jets with high  $P_T^{\text{jet}} > 30 \text{ GeV}$  in the region  $|\eta^{\text{jet}}| < 2.5$ . Among the isolated high  $P_T$  jets is required the presence at least of one tagged b-jet. This cut significantly suppresses  $WZ$  background and vanishes  $WH$ , the acceptance for them is  $(\sim 0.01 \div 0.05 \%)$ .**

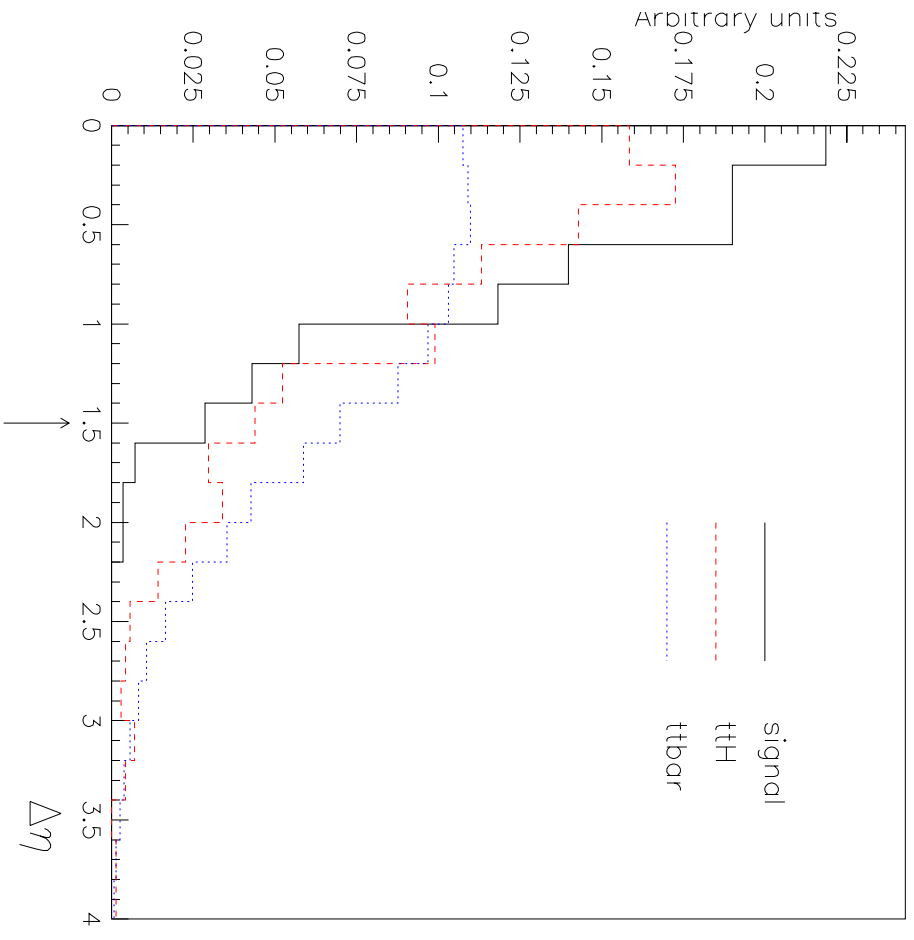
Cuts for $m_H=150, 160$ GeV	$t \rightarrow Hq$ Signal		Background Processes			
			$ttH$		WH	
	Nevt	Eff(%)	Nevt	Eff(%)	Nevt	Eff(%)
<b>Nevt gen.</b>	<b>10401</b> <b>8805</b>		<b>110000</b> <b>32800</b>		<b>300000</b> <b>15000</b>	
<b>Expected events</b>			<b>12950</b> <b>15020</b>		<b>430</b> <b>480</b>	
<b>Preselection</b>	<b>5012</b> <b>3943</b>	<b>48.00</b> <b>45.00</b>	<b>395</b> <b>475</b>	<b>3.05</b> <b>3.16</b>	<b>56</b> <b>66</b>	<b>13.07</b> <b>13.70</b>
<b><math>p_{T}^{\text{miss}} &gt; 45</math> GeV</b>						
<b><math>p_{T}^l &gt; 30</math> GeV</b>	<b>249</b>	<b>2.39</b>	<b>48</b>	<b>0.38</b>	<b>0</b>	<b><math>5.13 \times 10^{-2}</math></b>
<b><math>p_{T}^{\text{jet}} &gt; 30</math> GeV</b>	<b>218</b>	<b>2.48</b>	<b>76</b>	<b>0.51</b>	<b>0</b>	<b><math>5.33 \times 10^{-2}</math></b>
<b><math>N_{b\text{-jet}} \geq 1</math></b>						
<b><math>m_{ll} &lt; 80</math> GeV</b>	<b>186</b>	<b>1.79</b>	<b>22</b>	<b>0.17</b>	<b>0</b>	<b><math>4.03 \times 10^{-2}</math></b>
<b><math>\Delta\eta &lt; 1.5, \Delta\phi &lt; 1.0</math></b>	<b>170</b>	<b>1.93</b>	<b>35</b>	<b>0.24</b>	<b>0</b>	<b><math>4.67 \times 10^{-2}</math></b>
<b><math>t \rightarrow Hq, m_{ll} &lt; 110</math> GeV</b>	<b>87</b>	<b>0.84</b>	<b>6</b>	<b><math>5.1 \times 10^{-2}</math></b>	<b>0</b>	<b><math>7.30 \times 10^{-3}</math></b>
<b><math>p_{T}^j &gt; 30</math> GeV</b>	<b>60</b>	<b>0.68</b>	<b>12</b>	<b><math>8.5 \times 10^{-2}</math></b>	<b>0</b>	<b><math>6.70 \times 10^{-3}</math></b>
<b><math>t \rightarrow Wb, m_{lb} &lt; 140</math> GeV</b>	<b>78</b>	<b>0.75</b>	<b>6</b>	<b><math>4.6 \times 10^{-2}</math></b>	<b>0</b>	<b><math>6.70 \times 10^{-3}</math></b>
<b><math>p_{T}^{\text{bjet}} &gt; 40</math> GeV</b>	<b>59</b>	<b>0.67</b>	<b>11</b>	<b><math>7.6 \times 10^{-2}</math></b>	<b>0</b>	<b><math>5.10 \times 10^{-3}</math></b>

Cuts for $m_H=150, 160$ GeV	$t \rightarrow Hq$ Signal		Background Processes			
			$tt^-$		WZ	
	Nevt	Eff(%)	Nevt	Eff(%)	Nevt	Eff(%)
Nevt gen.	10401 8805		$4.0 \times 10^6$		$2 \times 10^5$	
Expected events			$3.8 \times 10^6$		39040	
Preselection	5012 3943	48.00 45.00	93611	2.41	3552	9.10
$p_{T}^{\text{miss}} > 45$ GeV $p_{T}^l > 30$ GeV $p_{T}^{\text{jet}} > 30$ GeV $N_{b\text{-jet}} \geq 1$	249 218	2.39 2.48	152	$3.91 \times 10^{-3}$	5	$1.28 \times 10^{-2}$
$m_{ll} < 80$ GeV $\Delta\eta < 1.5, \Delta\phi < 1.0$	186 170	1.79 1.93	38	$9.78 \times 10^{-4}$	0	$1.00 \times 10^{-3}$
$t \rightarrow Hq, m_{ll} < 110$ GeV $p_{T}^j > 30$ GeV	87 60	0.84 0.68	7	$1.80 \times 10^{-4}$	0	$5.00 \times 10^{-4}$
$t \rightarrow Wb, m_{lb} < 140$ GeV $p_{T}^{\text{bjet}} > 40$ GeV	78 59	0.75 0.67	3	$7.72 \times 10^{-5}$	0	$5.00 \times 10^{-4}$

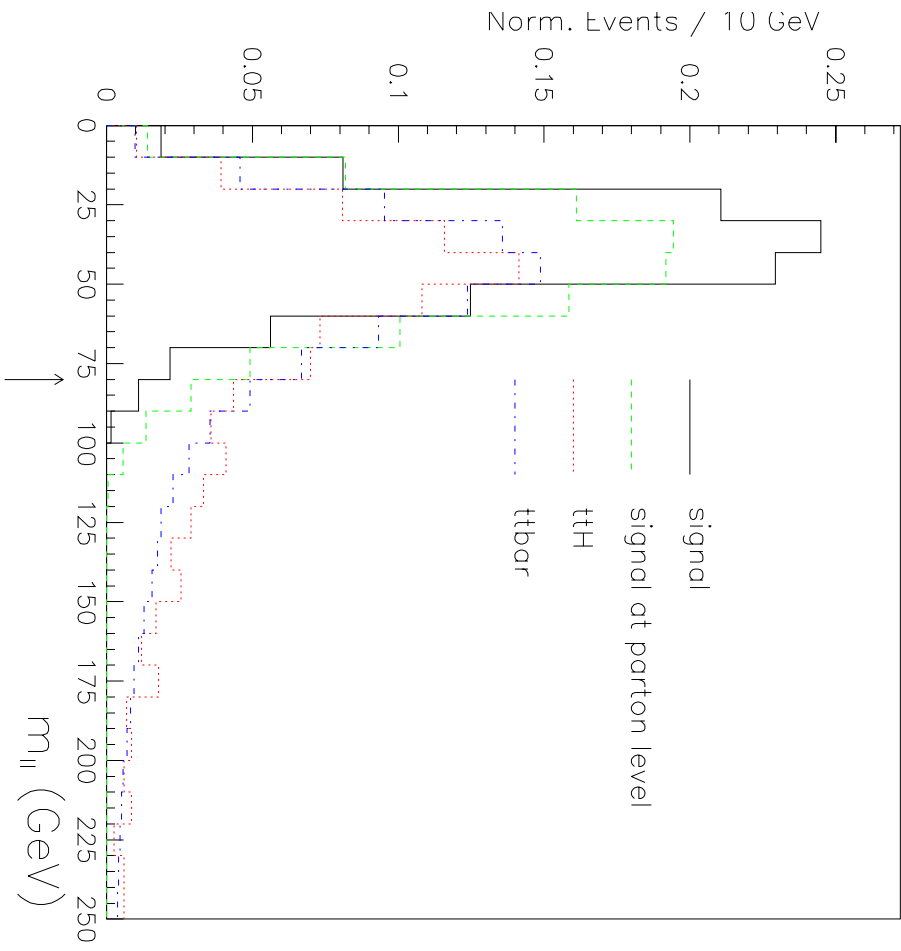


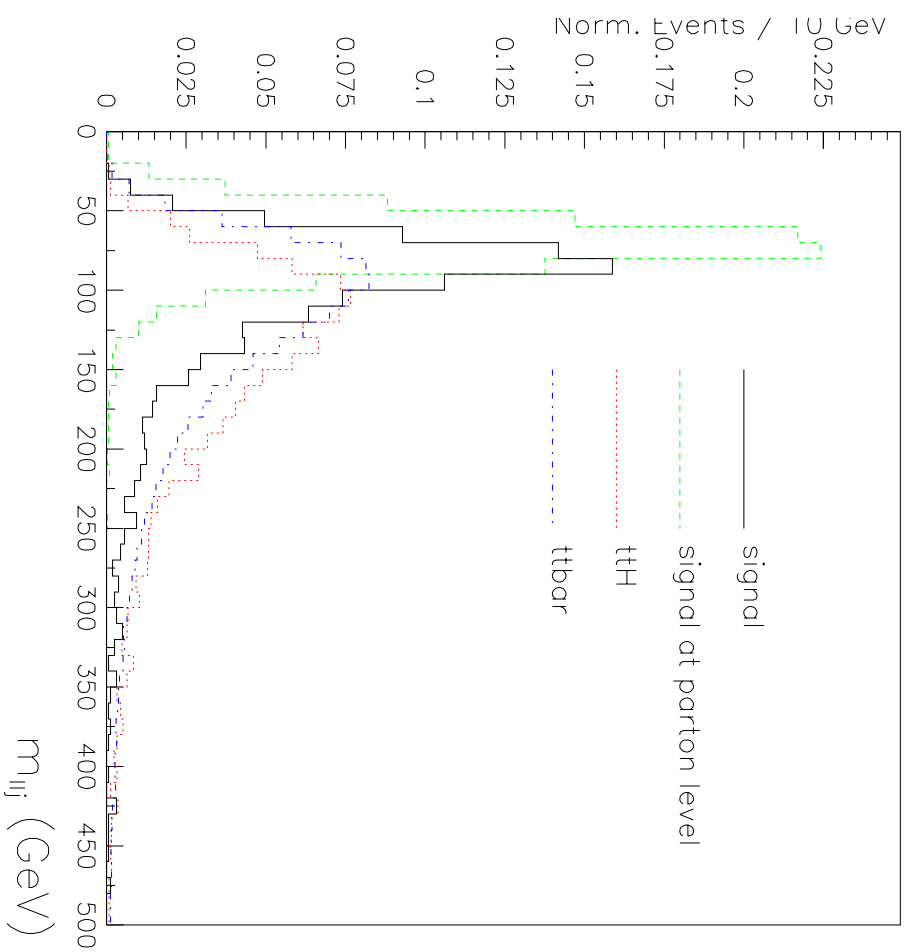






- For the dilepton (coming from **H** decay)  $m_{ll}$  mass reconstruction it has been required opposite charged lepton pairs with  $\Delta\phi < 1.0$  (the opening angle between the two leptons in the transverse plane, measured in rad.),  $\Delta\eta < 1.5$  (the absolute values of the pseudorapidity difference between the two leptons) and an invariant dilepton mass smaller than **80 GeV**. **These cuts sufficiently reduce  $t\bar{t}$  background** and **WZ background is vanished**.
- Then two kinematical cuts have been applied on:
  1. Invariant masses of  $m_{llj} < 110 \text{ GeV}$  with light jets  $P_T^j > 30 \text{ GeV}$  and
  2.  $m_{lb} < 140 \text{ GeV}$  with light jets  $P_T^{\text{bjet}} > 40 \text{ GeV}$
- It should be noted, that for  $m_{ll}$  and  $m_{llj}$  pairs the best combinations are defined as the closest values to the average numbers of **ll** and **llq** invariant mass distributions **at parton level**.





$$t\bar{t} \rightarrow HqWb \rightarrow WW^*qWb \rightarrow l\nu l\nu j, l\pm\nu b \quad (l = e, \mu)$$

- $S = 2. \times \sigma_{t\bar{t}} \times \text{Br}(t \rightarrow Hq) \times \text{Br}(t \rightarrow Wb) \times \text{Br}(H \rightarrow W^+ W^-) \times \text{Br}(W \rightarrow l\nu)^3 \times (\epsilon^l)^3 \times A^s \times L$

$\epsilon^l = 0.9$  – lepton identification efficiency

$A^s$  -- Signal Acceptance

$L$  -- Luminosity

- $S / \sqrt{(Bg \times (\epsilon^l)^3)} = 5.$

$Bg$  -- The sum of background events

- $\text{Br}(t \rightarrow Hq) = K \times 1 / A^s \times \sqrt{Bg}$

$$K = 5. / (2. \times \sigma_{t\bar{t}} \times \text{Br}(t \rightarrow Wb) \times \text{Br}(H \rightarrow W^+ W^-) \times \text{Br}(W \rightarrow l\nu)^3 \times L \times \sqrt{(\epsilon^l)^3})$$

$$K_{150} = 5. / (2. \times 833. \times 0.9982 \times 0.6852 \times 0.2163^3 \times 100000. \times$$

$$\times \sqrt{0.729})$$

$$K_{150} = 5.1 \times 10^{-6}$$

$$K_{160} = 3.8 \times 10^{-6}$$

## Summary of the results at $L = 100 \text{ fb}^{-1}$

Cuts	Sensitivity to $\text{Br}(t \rightarrow \text{H}q)$	
	$m_H = 150 \text{ GeV}$	$m_H = 160 \text{ GeV}$
$t \rightarrow \text{H}q,$ $m_{llj} < 110 \text{ GeV}$ $p_{j_T} > 30 \text{ GeV}$	$2.19 \times 10^{-3}$	$2.43 \times 10^{-3}$
$t \rightarrow \text{W}b,$ $m_{lb} < 140 \text{ GeV}$ $p_{b\text{jet}_T} > 40 \text{ GeV}$	$2.03 \times 10^{-3}$	$2.12 \times 10^{-3}$

The  $\text{Br}(t \rightarrow \text{H}q \rightarrow \text{W}W^*q)$  as low as  $2.0 \times 10^{-3}$  for  $m_H = 150 \text{ GeV}$  and as low as  $2.1 \times 10^{-3}$  for  $m_H = 160 \text{ GeV}$  could be discovered at the  $5\sigma$  level with an integrated luminosity of  $100 \text{ fb}^{-1}$ .

# Summary

- We have studied the **ATLAS sensitivity** to **FCNC top quark rare decay**  $t \rightarrow Hq$  ( $q=u,c$ ) with  $H \rightarrow WW^*$  at  $\sqrt{s} = 14$  **TeV** for an integrated luminosity of **100 fb<sup>-1</sup>**.
- The results demonstrate that, a branching ratio as low as  **$2.0 \times 10^{-3}$**  for  **$m_H = 150$  GeV** and as low as  **$2.1 \times 10^{-3}$**  for  **$m_H = 160$  GeV** could be discovered at the **5 $\sigma$**  level with an integrated luminosity of **100 fb<sup>-1</sup>**.



# Preliminary results

## for $tt^- \rightarrow HqWb \rightarrow bb^- j, l^\pm \nu b, (l=e,\mu)$ decays

- Study of the sensitivity of the **ATLAS experiment** to the **top quark rare decay via FCNC  $t \rightarrow Hq$  ( $q=u,c$ )** at  $\sqrt{s} = 14 \text{ TeV}$  in the decay mode of  $tt^- \rightarrow HqWb \rightarrow bb^- j, l^\pm \nu b, (l=e,\mu)$  is underway.
- The Standard Model backgrounds  $tt^-$ ,  $tt^- H$ ,  $W$ \_jets,  $WZ$ ,  $WH$  and  $Wbb^-$  have been analysed. The signal and backgrounds were generated via **PYTHIA 5.7** and **HERWIG** and simulated and analysed using **ATLFAST 2.51**. The preliminary result for the branching ratio  $t \rightarrow Hq \rightarrow bb^- j$  as low as  $4.5 \times 10^{-3}$  is obtained at  $5\sigma$  level with an integrated luminosity of  $100 \text{ fb}^{-1}$  for  $m_H = 115 \text{ GeV}$ .