

Lecture B2

INTERACTION of RADIATION with MATTER

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“Autumn Lectures” – Nuclear and Medical Physics

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OVERVIEW

SCATTERING

PHOTO EFFECT

COMPTON EFFECT

PAIR PRODUCTION

ATTENUATION

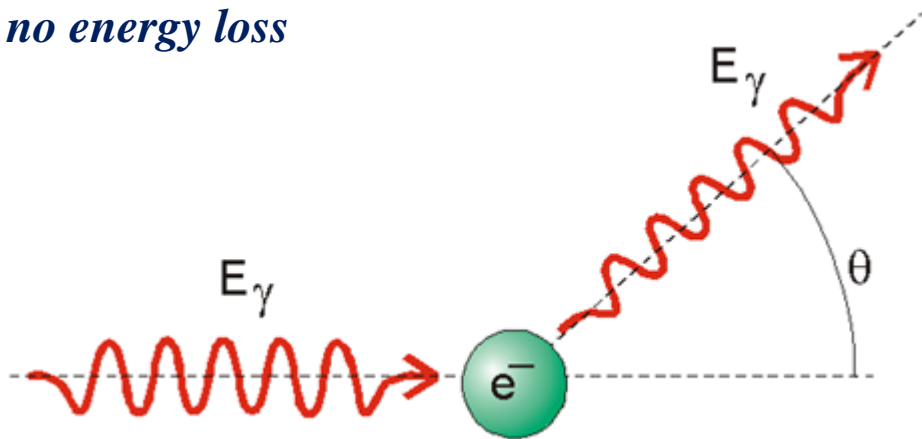
BREMSSTRAHLUNG

THOMSON SCATTERING

elastic scattering of el.-mag waves at a free charges = electron, ...

independent of wave length λ

no energy loss



application: plasma diagnosis, polarization of CMB, ...

$$\sigma_{Th, e} = \frac{8\pi}{3} \cdot r_e^2$$
$$\cong \frac{2}{3} \text{ barn}$$

$$\sigma_{Th, atom} = Z \cdot \sigma_{Th}$$

deviates from experiment

$$r_e = \alpha \cdot \frac{e^2}{4\pi\epsilon_0 m_e c^2}$$

$$= \alpha \cdot \frac{\hbar c}{m_e c^2}$$

$$= 2.82 \cdot 10^{-15} \text{ m}$$

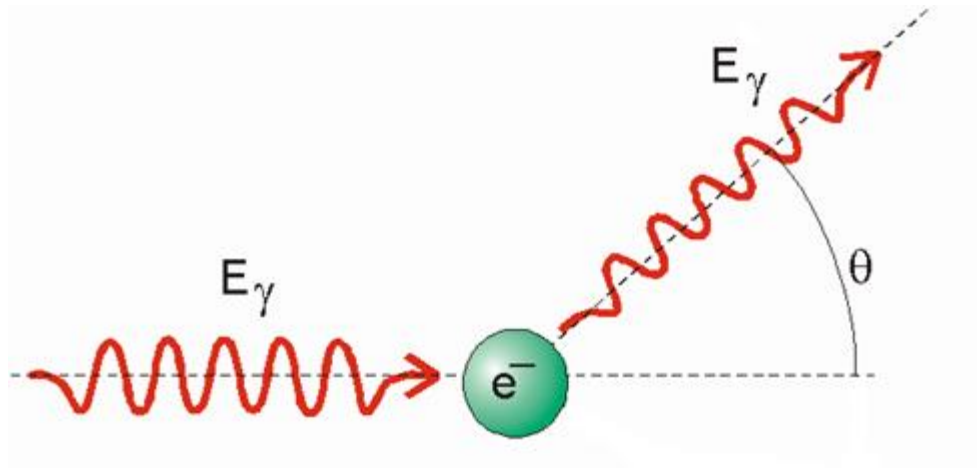
RAYLEIGH SCATTERING

elastic scattering of el.-mag waves at polarisable scattering centers = atoms, molecules

damped oscillation of „elastically“ bound electrons

eigen frequency ω_0

no energy loss



application: combustion diagnosis, holidays, ...

$$\sigma_R = \sigma_{Th} \cdot \frac{\omega^4}{(\omega^2 - \omega_0^2)^2} \cdot Z^2$$

ω_0 eigen frequency

$\omega \ll \omega_0$ makes the sky blue / sunset red

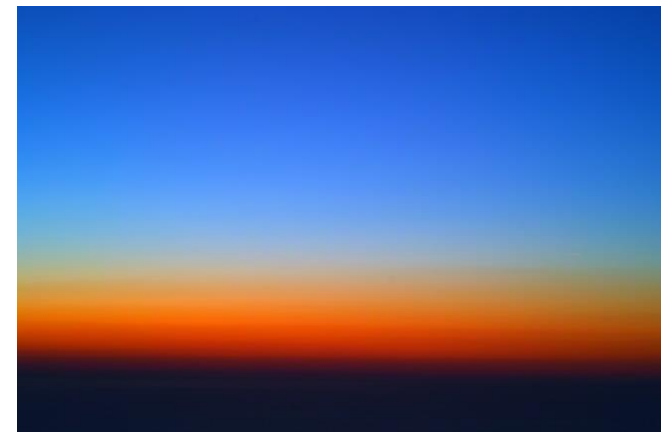


PHOTO EFFECT ✓

requires particle nature of „light“ Einstein 1905

$$\sigma_{2K} = \sigma_{Th} \cdot 4\sqrt{2}\alpha^4 \cdot Z^{4-5} \cdot \epsilon_\gamma^{-7/2}$$

$$\epsilon_\gamma = \frac{h\nu}{m_e c^2}$$

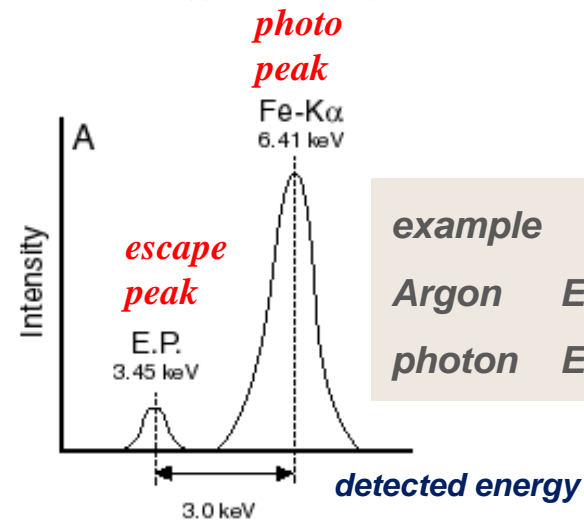
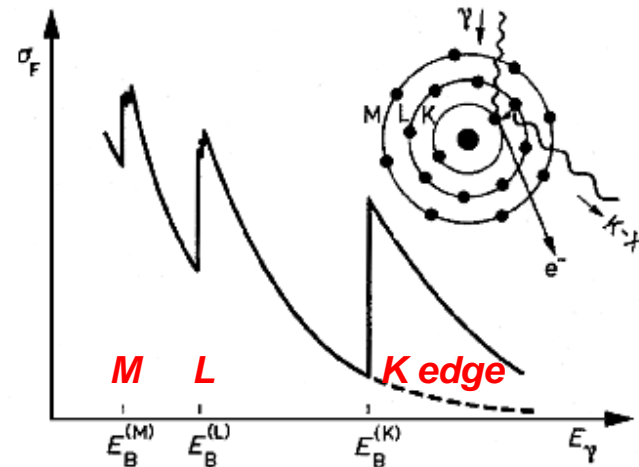
1. photon disappears
photo electron $E_e = E_{\text{photon}} - E_B$
2. refilling of hole in electron shell by
 - a) emission of photon or
 - b) Auger electron emission of loosely bound outer electron

$$E_{\text{Auger}} \cong E_B$$

detected energy E

photo peak $E = E_{\text{photon}}$
 $= E_e + E_B$

escape peak $E = E_{\text{photon}} - E_{K\alpha}$



example

Argon $E_{K\alpha} = 2.95 \text{ keV}$

photon $E_{\text{photon}} = 6.41 \text{ keV}$

COMPTON EFFECT ✓

proof of particle nature of „light“ Compton 1922

billiard with photons and „quasifree“ electrons

$$\sigma_C \approx \sigma_{Th} \cdot (1 - 2\varepsilon\gamma + \dots) \cdot Z \quad \varepsilon_\gamma \ll 1$$

$$\approx \sigma_{Th} \cdot \frac{3}{4} \cdot \left(\frac{1+2\ln\varepsilon\gamma}{2\varepsilon\gamma} + \dots \right) \cdot Z \quad \varepsilon_\gamma \gg 1$$

complicated QED calculation Klein&Nishina 1929

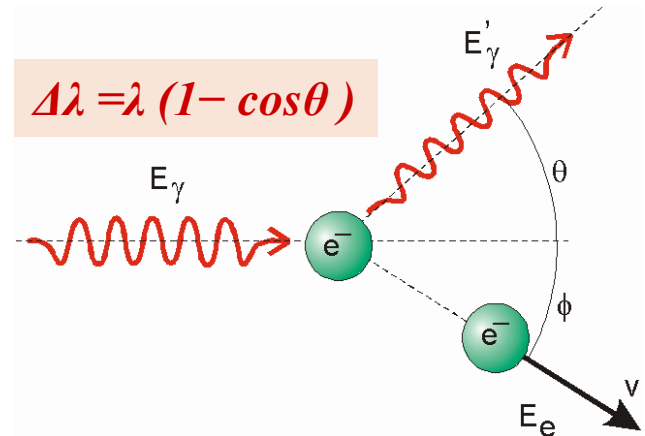
photon does not disappear

recoil electron

$$E_e = E_{\text{photon}} - E_{\text{photon}'}$$

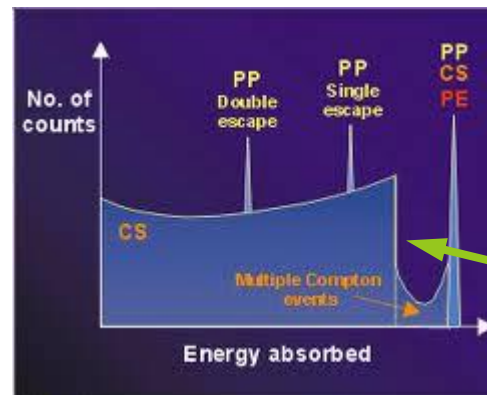
⇒

continuous spectrum



detected energy $E = E_e$

we neglect E_B of the electron and E_{recoil} of the nucleus because usually $E_B, E_{recoil} \ll E_e$



Compton edge = maximum energy transfer

PAIR PRODUCTION ✓

proof of mass-energy equivalence Blackett 1948

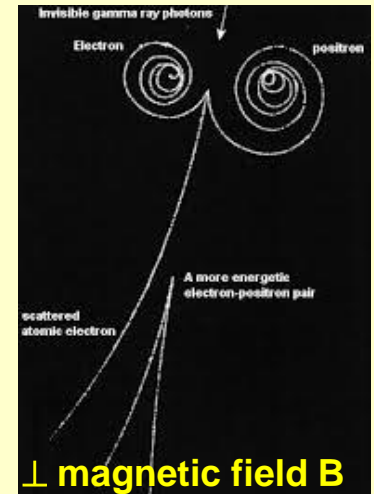
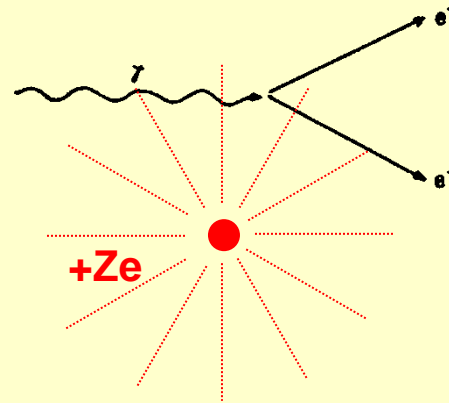
$$\sigma_{pair} \approx \sigma_{Th} \cdot Z^2 \cdot (\ln 2\epsilon\gamma + \dots) \quad \epsilon\gamma \gg 1$$

conversion of energy into matter

$$E_{\text{photon}} = h\nu > 2 m_{\text{electron, muon, pion, ...}}$$

a recoil partner (e.g. a nucleus) is needed to fulfil energy and momentum conservation

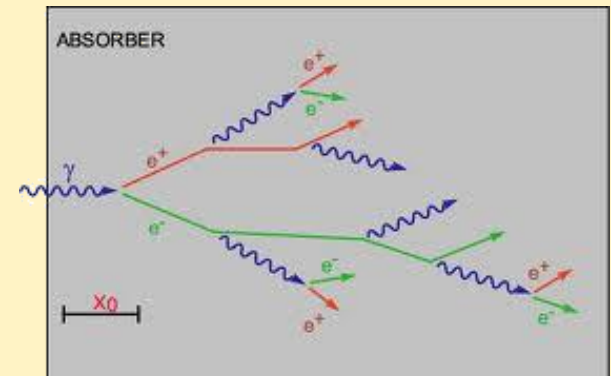
$$\begin{aligned} e^+e^- \text{ threshold: } m_{\text{recoil}} &= \infty & h\nu &= 2 m_e c^2 \\ & & &= m_e &= 4 m_e c^2 \end{aligned}$$



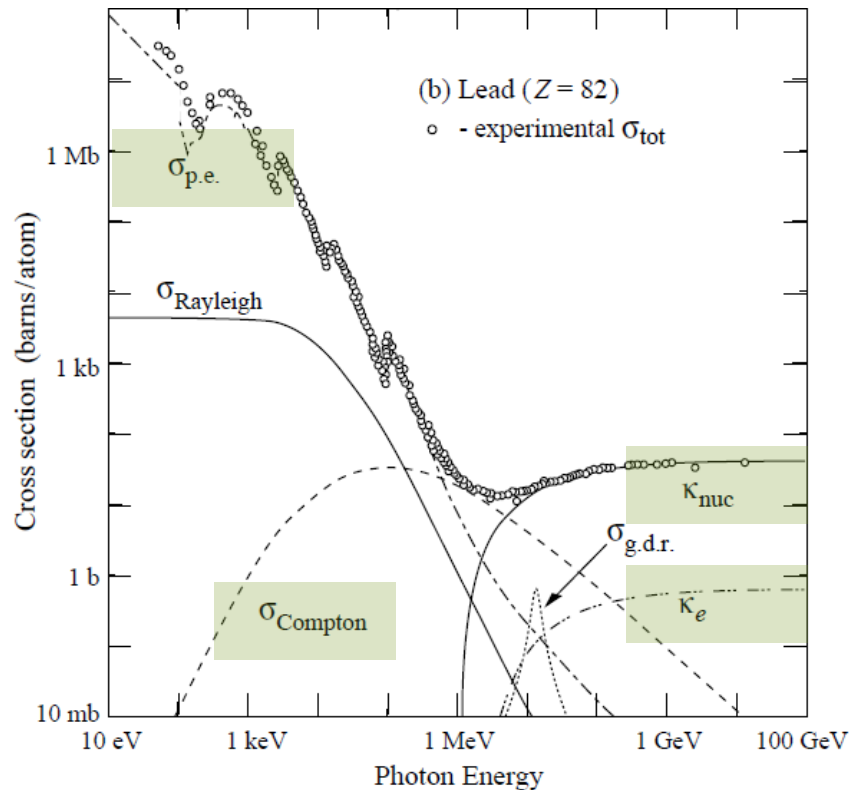
el.-mag shower

$e^+ e^- \gamma$ - cascade
pair production and bremsstrahlung alternate
shower may start with photon or electron

radiation length X_0
characteristic material dependent constant
depth, where about 2/3 of the incident energy is converted

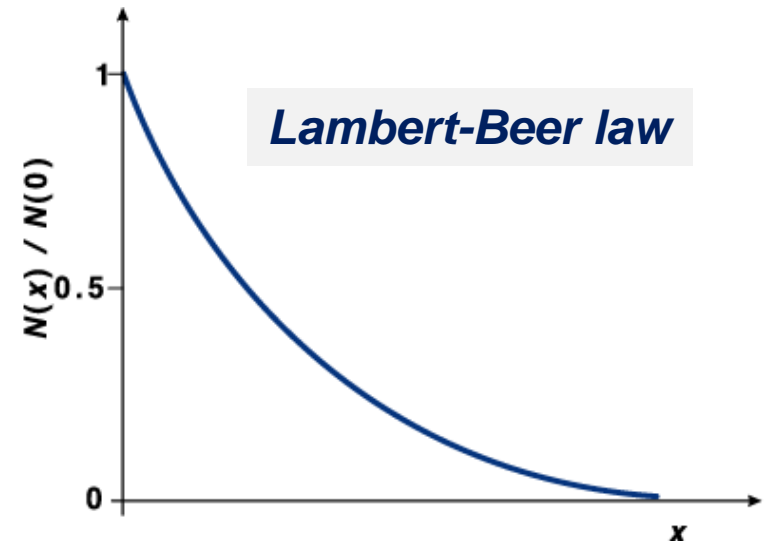
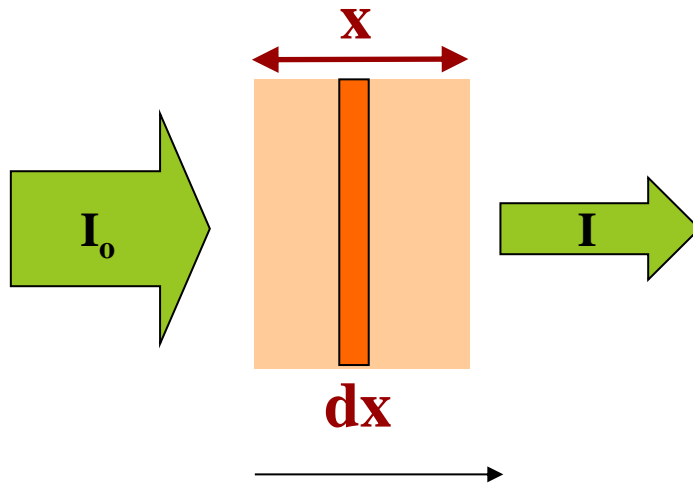


CROSS SECTIONS SUMMARY



- $\sigma_{\text{p.e.}}$ = Atomic photoelectric effect (electron ejection, photon absorption)
- σ_{Rayleigh} = Rayleigh (coherent) scattering—atom neither ionized nor excited
- σ_{Compton} = Incoherent scattering (Compton scattering off an electron)
- κ_{nuc} = Pair production, nuclear field
- κ_e = Pair production, electron field
- $\sigma_{\text{g.d.r.}}$ = Photonuclear interactions,

ATTENUATION



$$I(x) = I_0 e^{-\mu(h\nu)x} \quad \text{intensity after layer thickness } x$$

$$\frac{I_T(x)}{I_0} = e^{-\mu(h\nu)x} \quad \text{fraction of transmission}$$

$$\frac{I_A(x)}{I_0} = 1 - e^{-\mu(h\nu)x} \quad \text{fraction of absorption}$$

sum of linear attenuation coeff.

$$\mu(h\nu) = \sum_i \mu_i(h\nu)$$

$$\mu_i = \rho \cdot \frac{N_A}{A} \cdot \sigma_i(h\nu)$$

BREMSSTRAHLUNG

accelerated charged particles radiate Hertz 1886

electromagnetic waves

$$\sigma_b \approx \sigma_{Th} \cdot Z^2 \cdot [\text{energy dependent}]$$

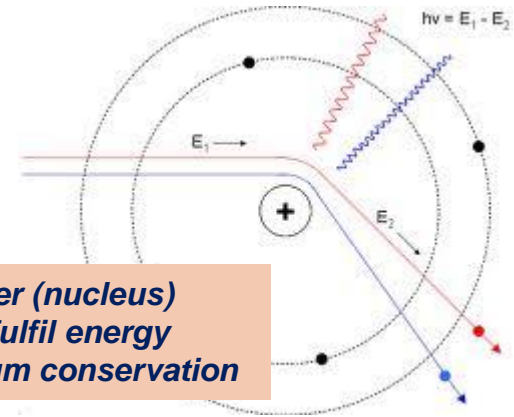
bending force by Coulomb potential

force \Leftrightarrow acceleration

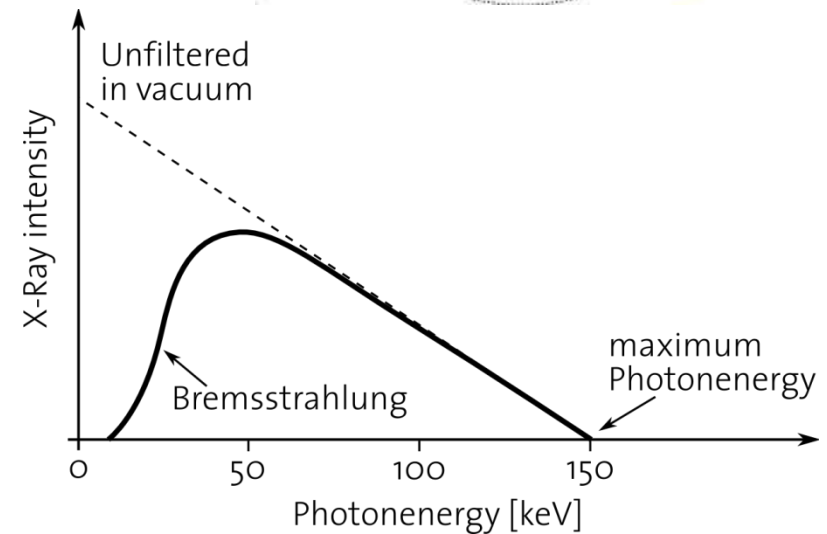
$$F_{\text{Coulomb}} = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q_{\text{particle}} \cdot Q_{\text{nucleus}}}{r^2}$$
$$= m \cdot \ddot{r}$$

any distance r

\Rightarrow *continuous spectrum*



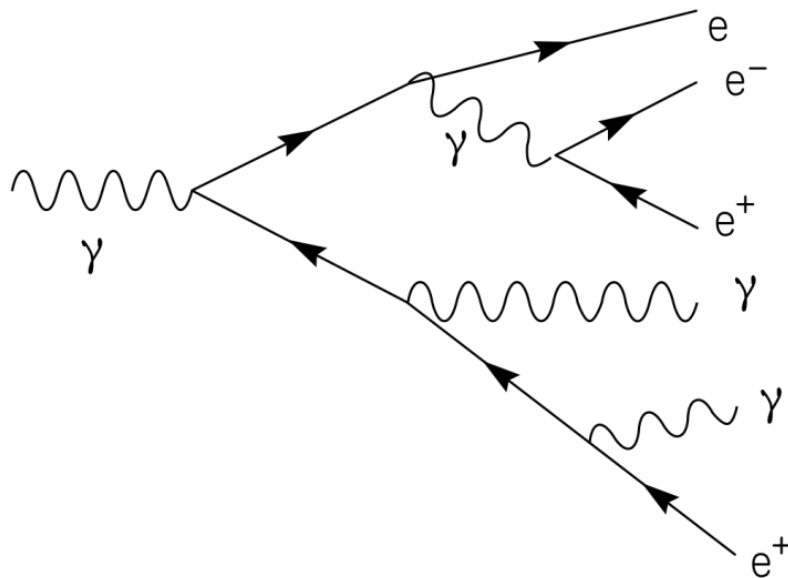
*a recoil partner (nucleus)
is needed to fulfil energy
and momentum conservation*



EL.-MAG. SHOWER

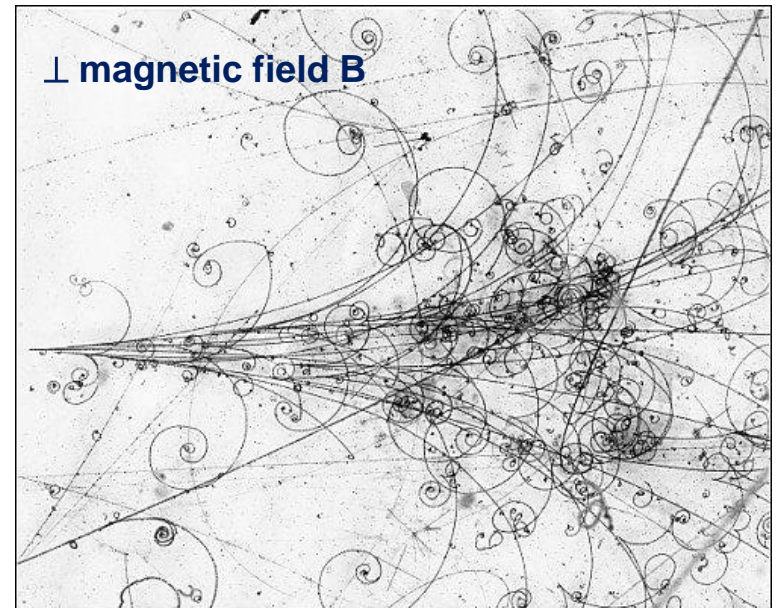
alternating pair production & bremsstrahlung

initial particle of minor importance for large energies



characteristic quantity of **absorber**

radiation length X_0



$$E_\gamma = E_{initial} \cdot e^{-(x/X_0)}$$